

An introduction to Topological Quantum Computing

Shawn Xingshan Cui
Dept. of Mathematics
Dept. of Physics and Astronomy
Purdue University

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①

(physics)
topological phase of matter



topological
quantum computing

②

(algebra)
higher category theory



③

(topology)
topological quantum field theory



Part III: Outline

- TQC Setup
- Universality of anyons
- Quantum invariants from anyons

TQC Setup

Fix an anyon system or UMTC $\{L, N_{ab}^c, F_{d,ef}^{abc}, R_c^{ba}\}$

★ Qubit spaces and computational bases

$V_c^{a,\dots,a} \equiv V_b^{a^{\otimes n}}$: space of physical processes of splitting b into n anyons of type a ; space of n type- a anyons with total type b

$$V_b^{a^{\otimes n}} = \begin{array}{c} \text{---} \bullet \text{---} \bullet \text{---} \bullet \text{---} \cdots \text{---} \bullet \text{---} \\ \text{\textit{a}} \quad \text{\textit{a}} \quad \text{\textit{a}} \quad \cdots \quad \text{\textit{a}} \\ \text{---} \end{array}$$

b

- a must be a non-Abelian anyon

$$\dim V_b^{a^{\otimes n}} \sim (\dim_q a)^n$$

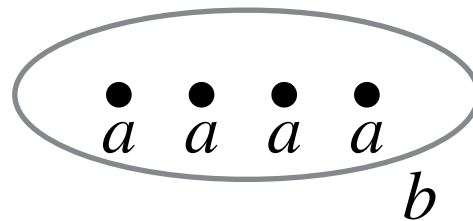
- In general, $V_b^{a^{\otimes n}}$ does not have a tensor product structure. To encode k qubits, choose n large enough so that

$$(\mathbb{C}^2)^k \hookrightarrow V_b^{a^{\otimes n}}$$

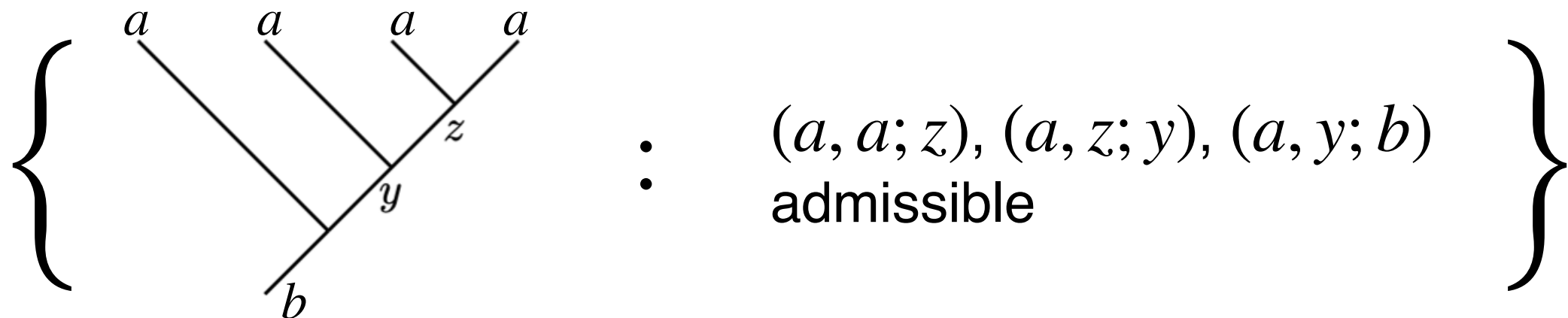
Practically: want a systematic way to choose the embedding

- Choose the 'computing' anyon a , 'total-type' anyon b , and a fixed integer m , such that $V_b^{a^{\otimes m}}$ is the desired qubit/qudit.

e.g. $m = 4$

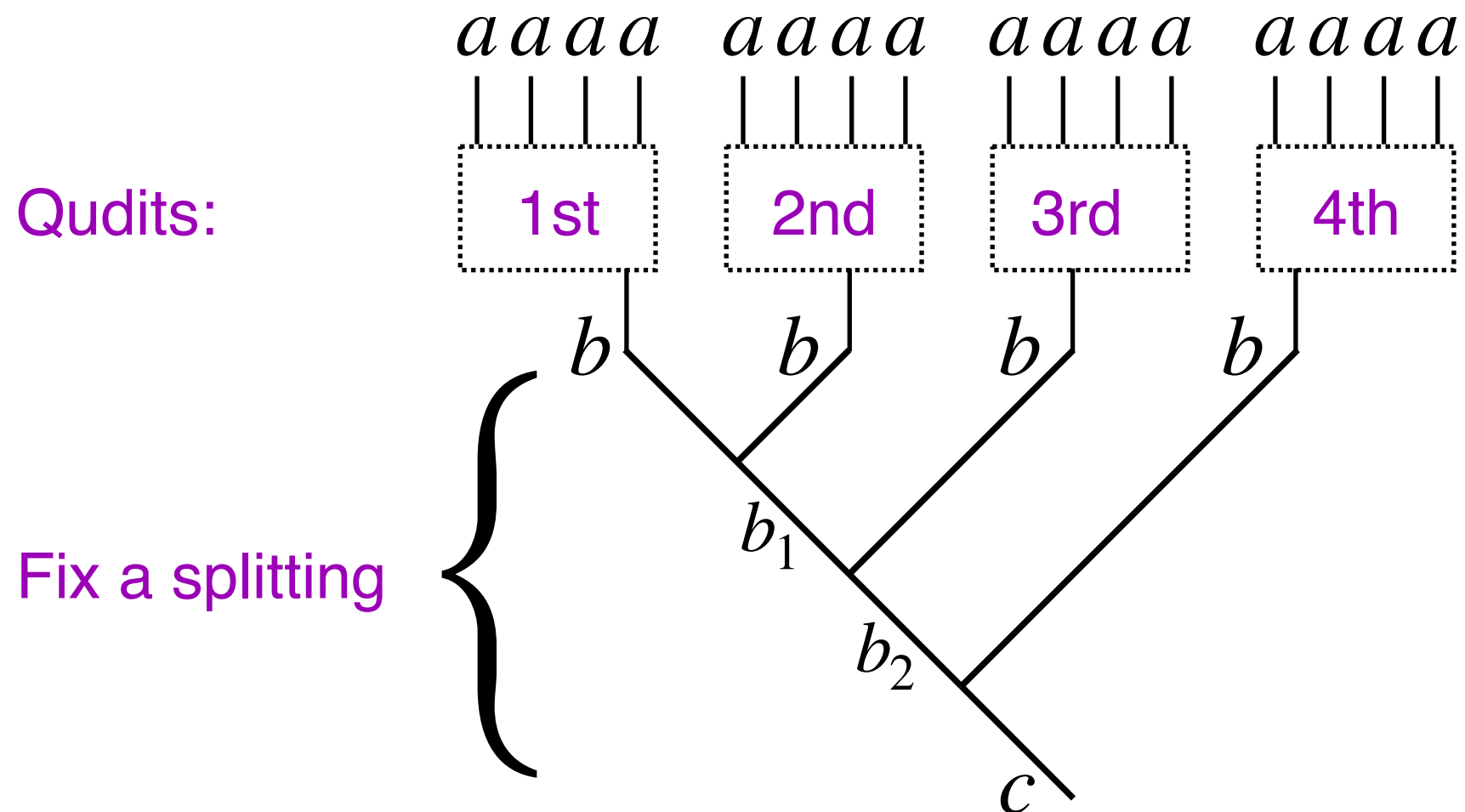


Choose a splitting-tree basis as the computational basis:

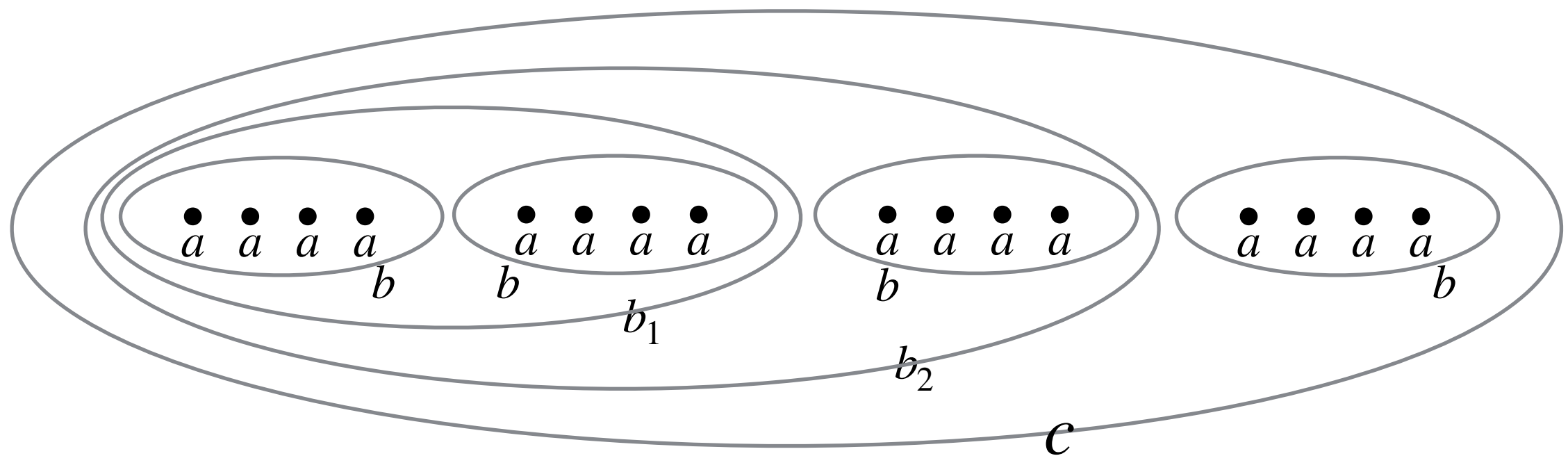


- To encode k qudits, choose anyon c such that $V_c^{b^{\otimes k}} \neq 0$

$$\underbrace{V_b^{a^{\otimes m}} \otimes \dots \otimes V_b^{a^{\otimes m}}}_k \hookrightarrow V_c^{a^{\otimes mk}}$$



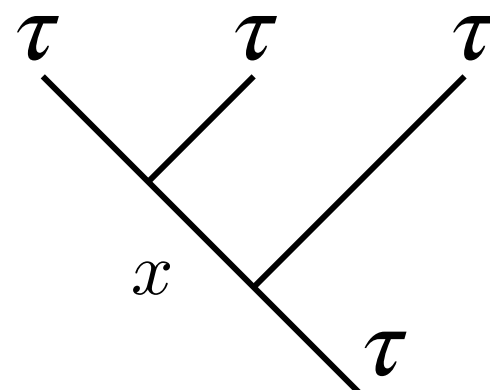
2D picture:



Example: Fib

$$\tau \otimes \tau = \mathbf{1} \oplus \tau$$

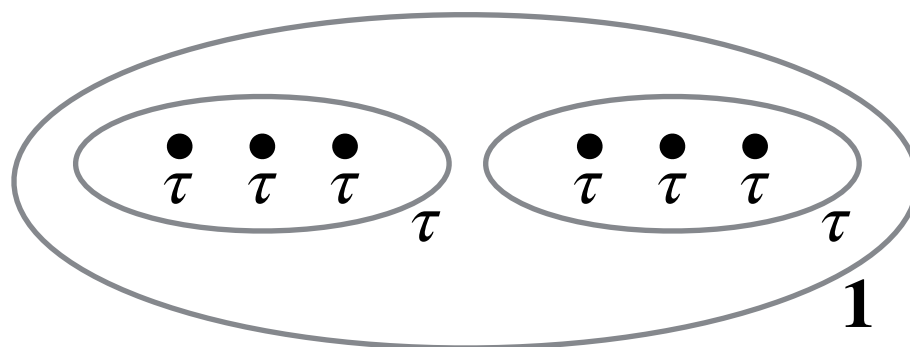
$V_{\tau}^{\tau\tau\tau}$ has a basis:



$$x = \mathbf{1}, \tau$$

Hence $V_{\tau}^{\tau\tau\tau}$ is a qubit.

Two qubits: $V_{\tau}^{\tau\otimes 3} \otimes V_{\tau}^{\tau\otimes 3} \hookrightarrow V_{\mathbf{1}}^{\tau\otimes 6}$

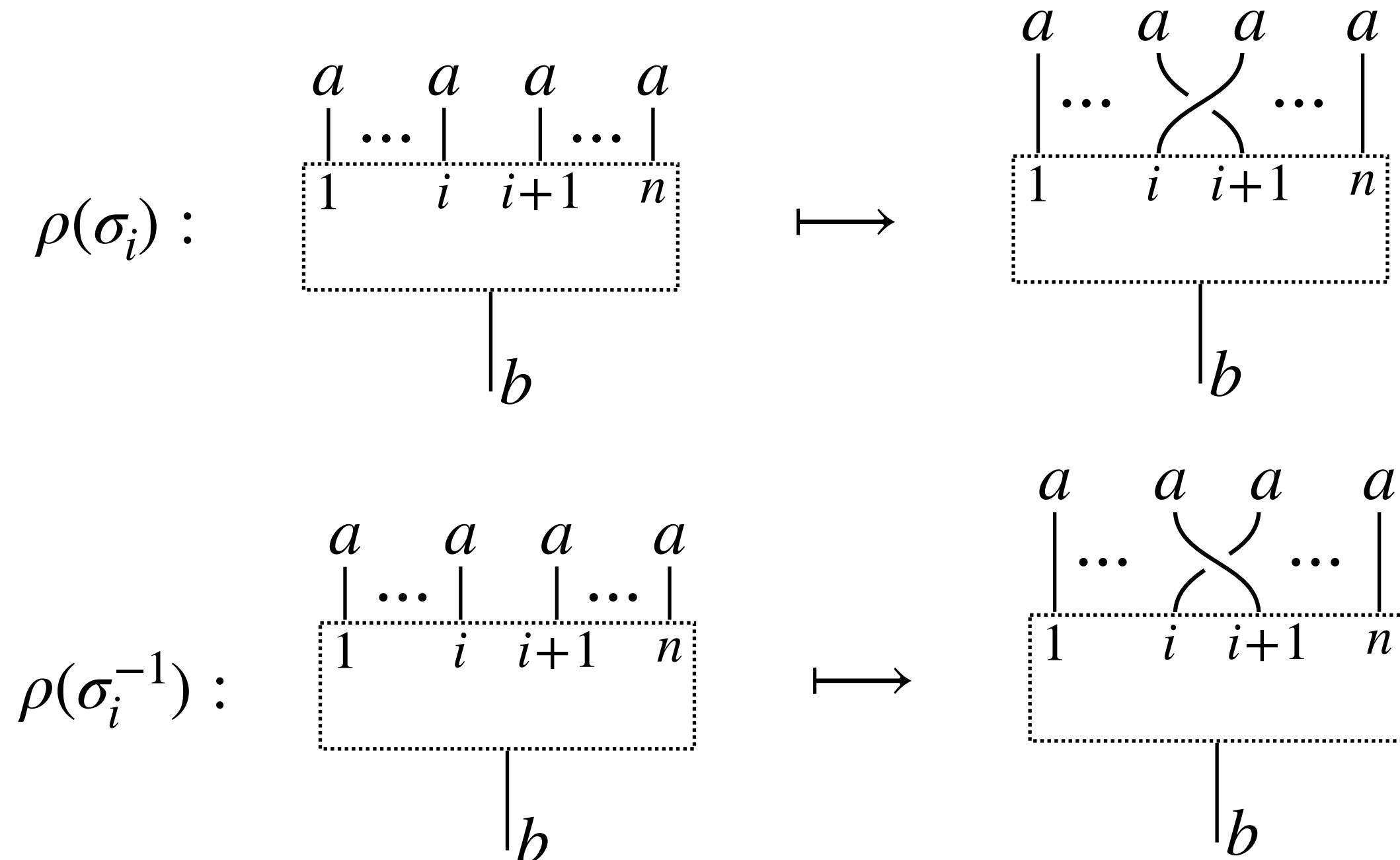


Note $\dim V_{\tau}^{\tau\otimes 6} = F_5 = 5$

★ Quantum gates from braiding

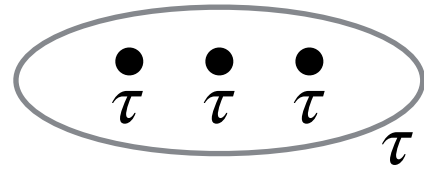
Recall: braiding induces unitary representation of braid groups

$$\rho \equiv \rho_b^{a,n} : B_n \longrightarrow U(V_b^{a^{\otimes n}})$$



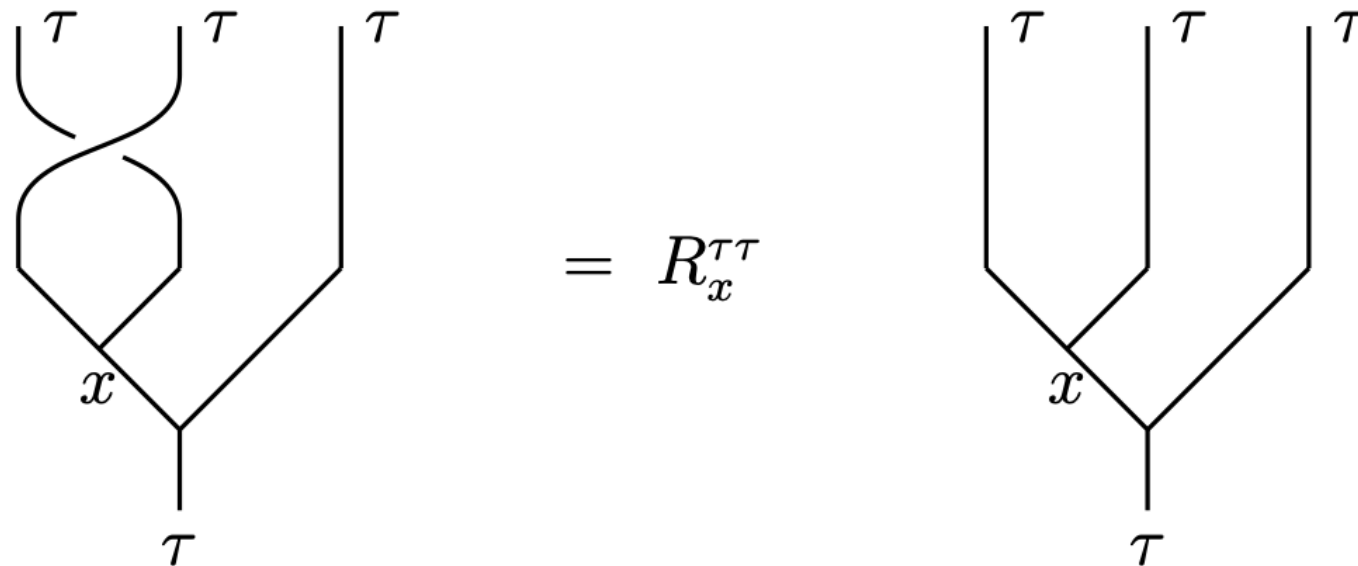
Example: Fib

- 1-qubit gates

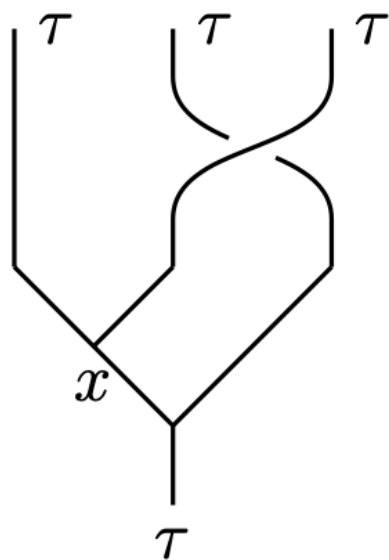


$$\rho : B_3 = \langle \sigma_1, \sigma_2 \rangle \longrightarrow U(V_\tau^{\tau \otimes 3})$$

$$x = \mathbf{1}, \tau$$

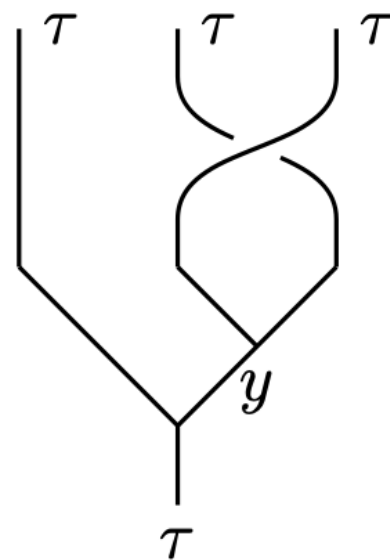


$$\rho(\sigma_1) = \begin{pmatrix} R_1^{\tau\tau} & 0 \\ 0 & R_\tau^{\tau\tau} \end{pmatrix} = \begin{pmatrix} e^{-4\pi i/5} & 0 \\ 0 & e^{3\pi i/5} \end{pmatrix}$$

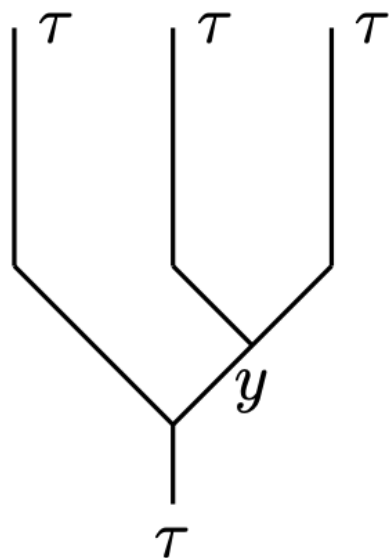


=

$$\sum_{y=\mathbf{1}, \tau} F_{\tau;yx}^{\tau\tau\tau}$$

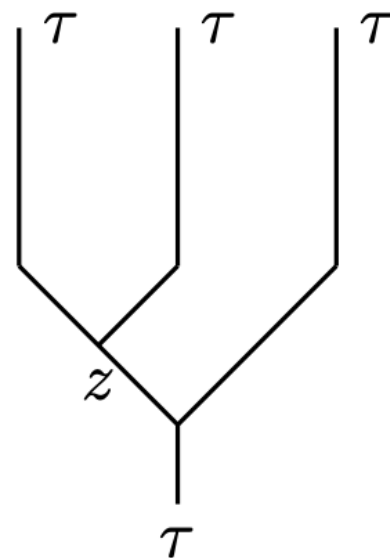


$$= \sum_{y=\mathbf{1}, \tau} F_{\tau;yx}^{\tau\tau\tau} R_y^{\tau\tau}$$



=

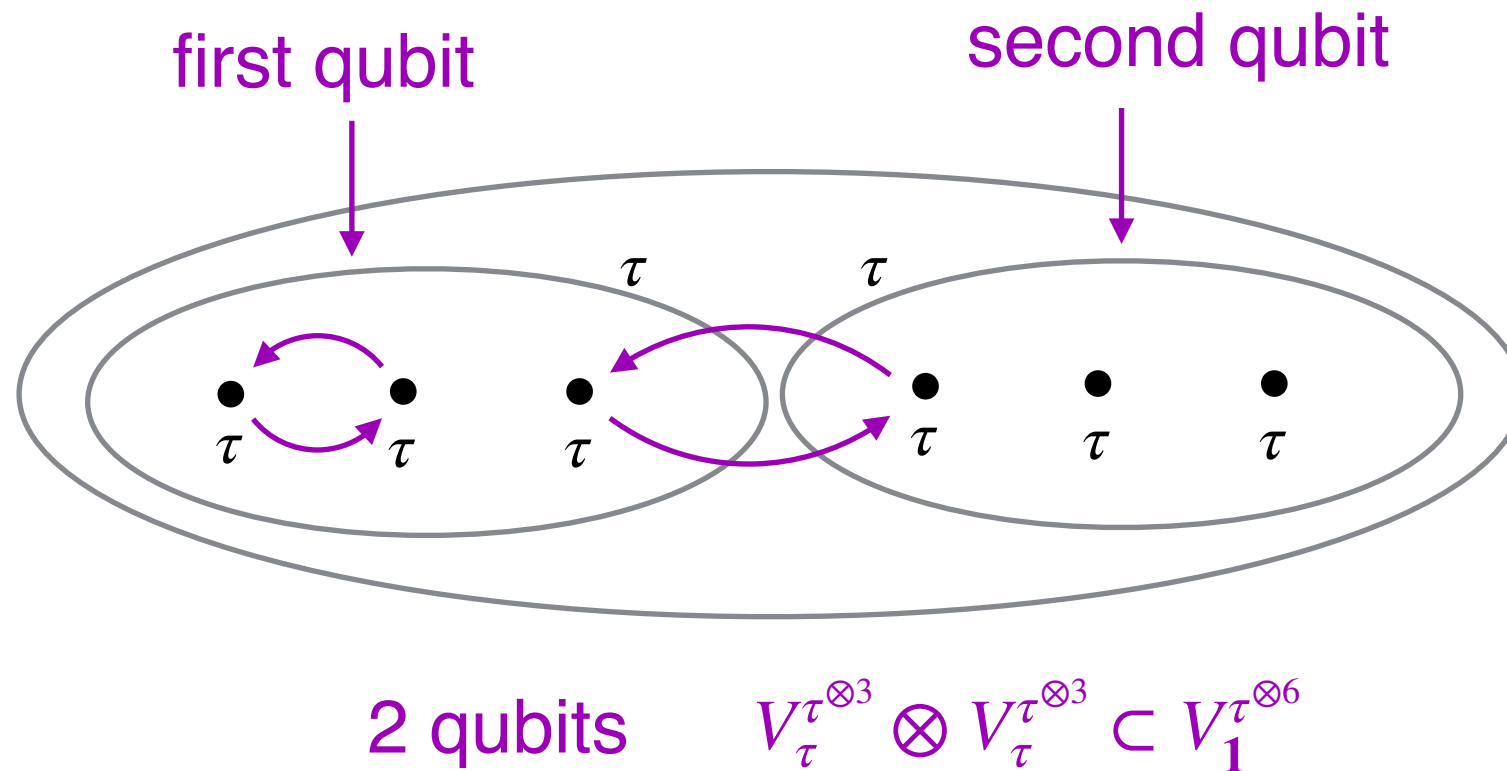
$$\sum_{y,z=\mathbf{1}, \tau} F_{\tau;yx}^{\tau\tau\tau} R_y^{\tau\tau} (F_{\tau}^{\tau\tau\tau})_{zy}^{-1}$$



$$\rho(\sigma_2)_{zx} = \sum_{y=\mathbf{1}, \tau} F_{\tau;yx}^{\tau\tau\tau} R_y^{\tau\tau} (F_{\tau}^{\tau\tau\tau})_{zy}^{-1}$$

$$\rho(\sigma_2) = (F_{\tau}^{\tau\tau\tau})^{\dagger} \rho(\sigma_1) F_{\tau}^{\tau\tau\tau}$$

- Multi-qubit gates are computed similarly



A general braiding acts on $V_1^{\tau^{\otimes 6}}$, not necessarily preserve $V_{\tau}^{\tau^{\otimes 3}} \otimes V_{\tau}^{\tau^{\otimes 3}}$

Conjecture

In the above encoding, there is no leakage-free entangling braiding gates.

- A braiding on $V_1^{\tau^{\otimes 6}}$ that preserves $V_{\tau}^{\tau^{\otimes 3}} \otimes V_{\tau}^{\tau^{\otimes 3}}$, and
- On $V_{\tau}^{\tau^{\otimes 3}} \otimes V_{\tau}^{\tau^{\otimes 3}}$, is not of the form $A \otimes B$ or $\text{SWAP} \circ (A \otimes B)$

Universality of anyons

An anyon a is **braiding universal** if for some anyon b and for all N large enough, the image of $\text{Im } \rho_b^{a,N}$ of the braid group is dense in $U(V_b^{a^{\otimes N}})$.

Theorem (Freedman-Larsen-Wang, 2002)

The τ anyon in Fib is braiding universal.

Theorem (Freedman-Larsen-Wang, 2002)

For the anyon theory $\text{SU}(2)_k$ arising from the quantum group $U_q(\mathfrak{sl}_2)$, $q = e^{2\pi i/(k+2)}$, if k is not 0, 1, 2, 4, then the non-Abelian anyons are braiding universal. (Fib $\sim \text{SU}(2)_3$)

Property F Conjecture (Naidu-Rowell, 2011)

An anyon has finite braid group image if and only if its quantum dimension squared is an integer. **(True for all known cases)**

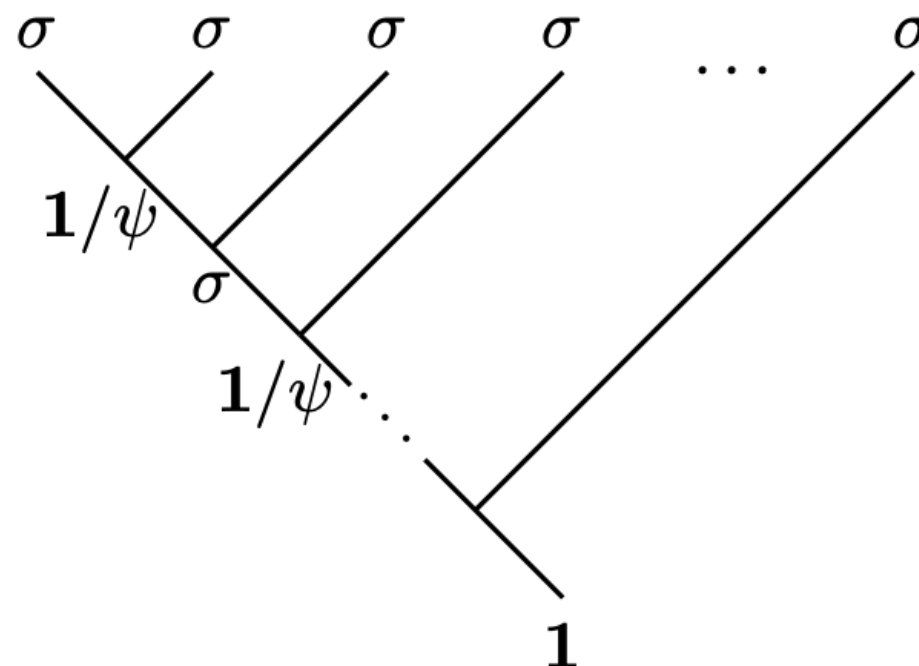
Example: Ising = $SU(2)_2$

- Label set: $L = \{\mathbf{1}, \psi, \sigma\}$; $\bar{a} = a$.
- Fusion rules: $\psi \otimes \psi = \mathbf{1}$, $\psi \otimes \sigma = \sigma$, $\sigma \otimes \sigma = \mathbf{1} \oplus \psi$.
- F -matrices:

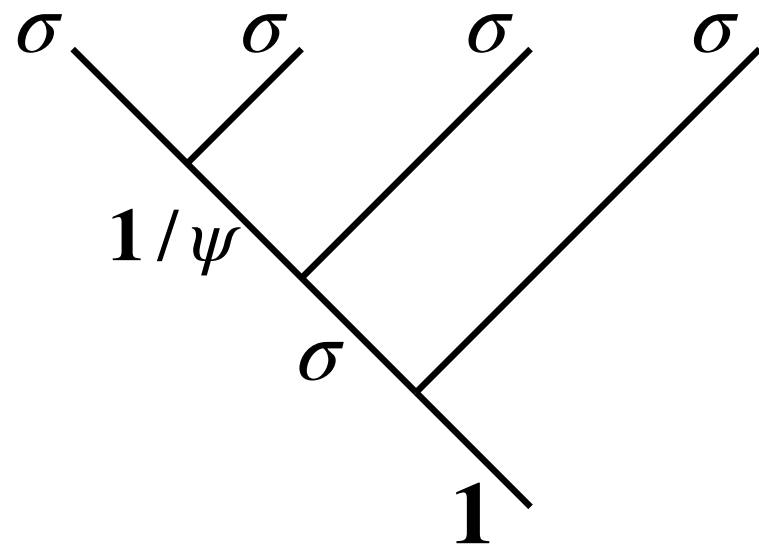
$$F_{\sigma}^{\sigma\sigma\sigma} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad F_{\sigma}^{\psi\sigma\psi} = F_{\psi}^{\sigma\psi\sigma} = (-1).$$

- R -symbols: $R_1^{\psi\psi} = -1$, $R_{\sigma}^{\psi\sigma} = R_{\sigma}^{\sigma\psi} = -i$, $R_1^{\sigma\sigma} = e^{-\pi i/8}$, $R_{\psi}^{\sigma\sigma} = e^{3\pi i/8}$.

$\dim V_1^{\sigma^{\otimes 2n}} = 2^{n-1}$, a rare case!



Qubit $V_1^{\sigma\sigma\sigma\sigma}$:



$$\rho(\sigma_1) = \rho(\sigma_3) = e^{-\frac{\pi i}{8}} \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix},$$

$$\rho(\sigma_2) = e^{-\frac{\pi i}{8}} \begin{pmatrix} \frac{1-i}{2} & \frac{1+i}{2} \\ \frac{1+i}{2} & \frac{1-i}{2} \end{pmatrix}$$

Fact: $\text{Im } \rho_1^{\sigma, 2n}$ is the Clifford group generated by $\{\text{CNOT}, H, P, X\}$

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \quad P = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix} \quad X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \text{CNOT} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

Theorem (Gottesman-Knill)

Quantum gates in the Clifford group can be simulated efficiently on a classical computer.

A quantum computer built from the Ising anyons have no quantum advantage.

Example: $SU(2)_4$

Anyon types $L = \{0, 1, 2, 3, 4\}$

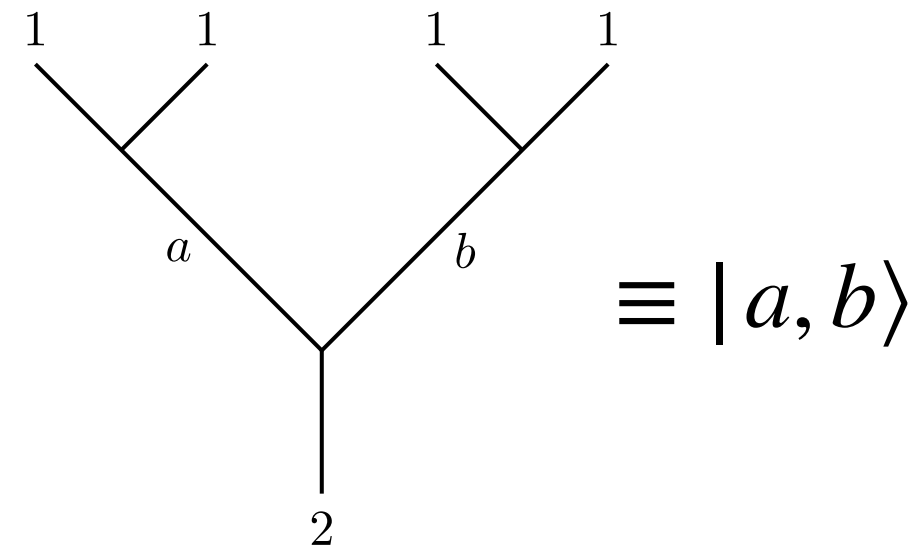
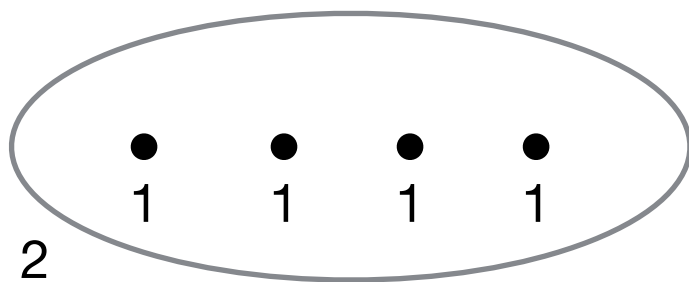
vacuum

Partial fusion rules:

$$1 \otimes 1 = 0 \oplus 2$$

$$2 \otimes 2 = 0 \oplus 2 \oplus 4$$

- qutrit model

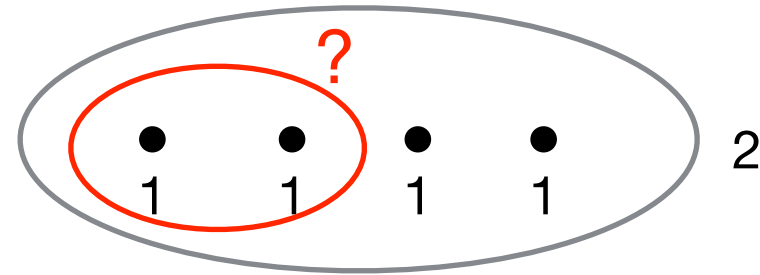


$$(a, b) = (0, 2), (2, 0), (2, 2)$$

The braid group has finite image — not braiding universal

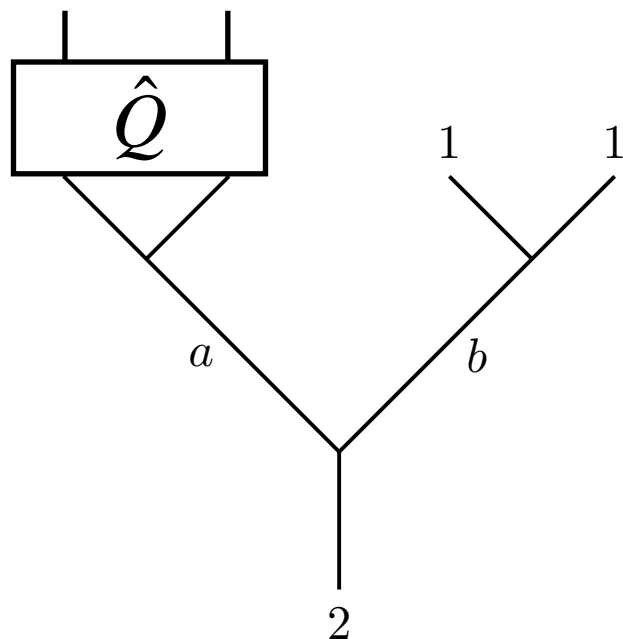
- Measurement of total type

$$1 \otimes 1 = 0 \oplus 2$$



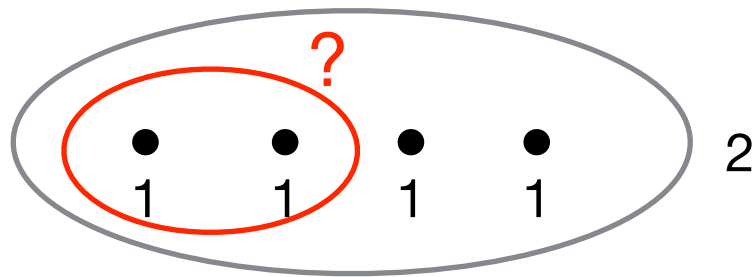
For some $\lambda_0 \neq \lambda_2 \in \mathbb{R}$

$$\hat{Q} := \lambda_0 \begin{array}{c} 1 \quad 1 \\ \diagdown \quad \diagup \\ \\ \diagup \quad \diagdown \\ 1 \quad 1 \end{array} 0 + \lambda_2 \begin{array}{c} 1 \quad 1 \\ \diagdown \quad \diagup \\ \\ \diagup \quad \diagdown \\ 1 \quad 1 \end{array} 2$$



$$= \lambda_a \begin{array}{c} 1 \quad 1 \quad 1 \quad 1 \\ \diagdown \quad \diagup \quad \diagdown \quad \diagup \\ \\ \diagup \quad \diagdown \\ 2 \end{array}$$

$$(a, b) = (0, 2), (2, 0), (2, 2)$$



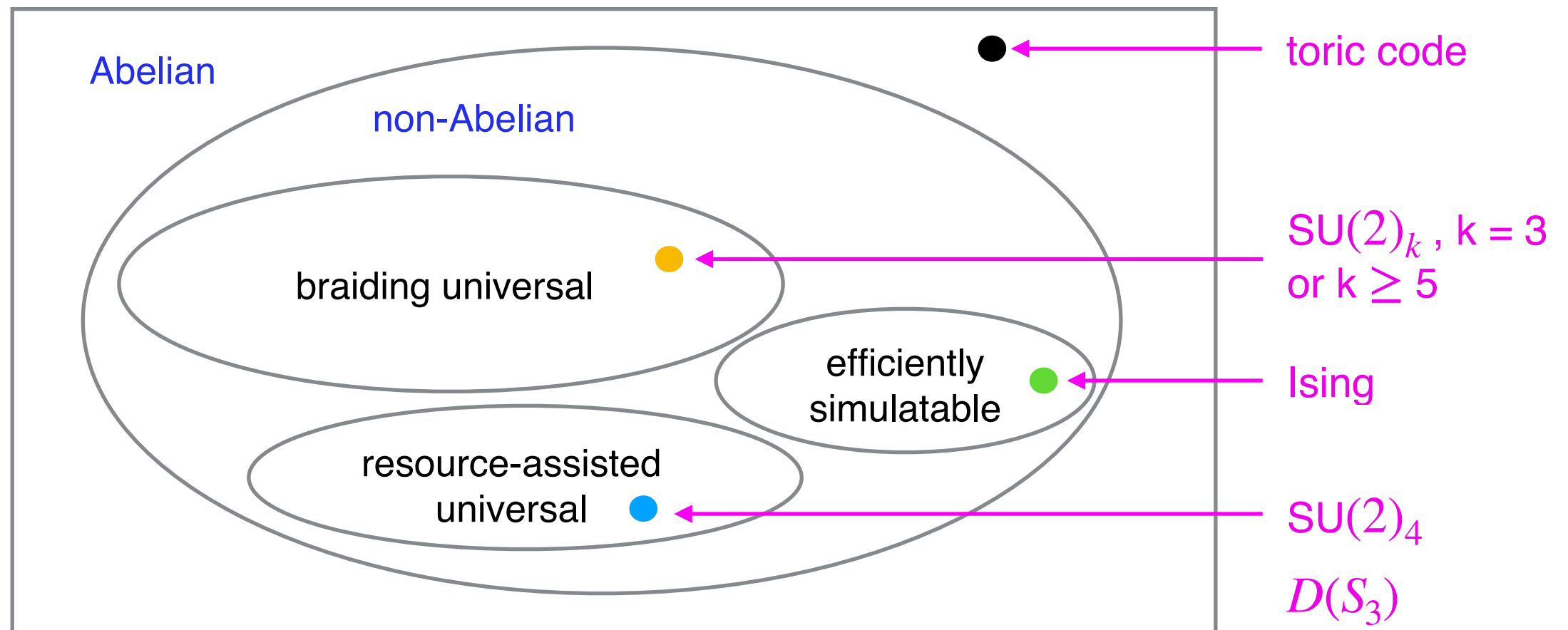
Effect of measuring the total type of the two type-1 anyons, i.e. $\hat{Q}$:

$$\begin{array}{lcl}
 \alpha |0,2\rangle + \beta |2,0\rangle + \gamma |2,2\rangle & \begin{array}{l} \nearrow |0,2\rangle \\ \searrow \frac{\beta |2,0\rangle + \gamma |2,2\rangle}{|\beta|^2 + |\gamma|^2} \end{array} & \begin{array}{l} \text{Prob} = |\alpha|^2 \\ \text{Prob} = |\beta|^2 + |\gamma|^2 \end{array} \\
 |\alpha|^2 + |\beta|^2 + |\gamma|^2 = 1 & &
 \end{array}$$

Theorem (Cui - Wang, '14)

In $SU(2)_4$ with the above encoding, braiding supplemented with the measurement of the total type of two anyons, is universal —

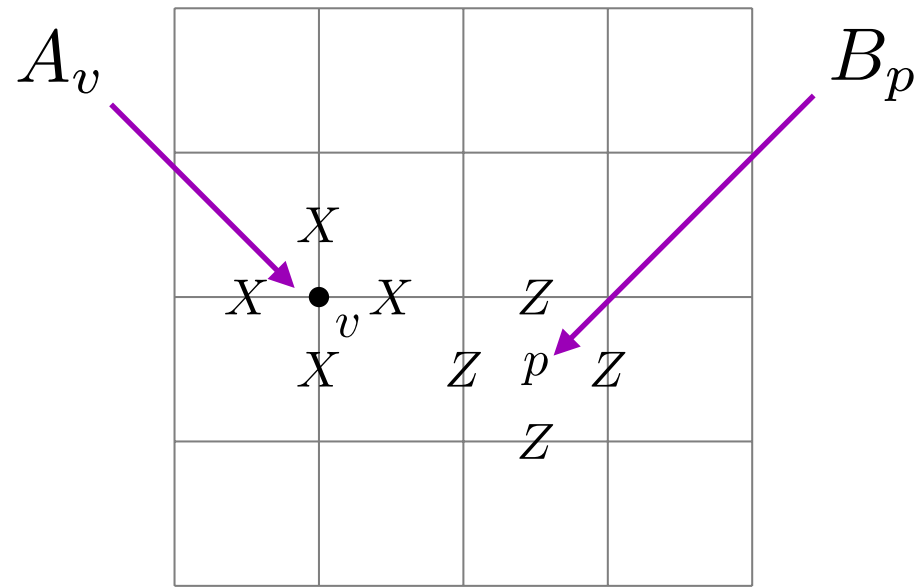
Resource-Assisted Universality



Problem

How to characterize each class of anyons?

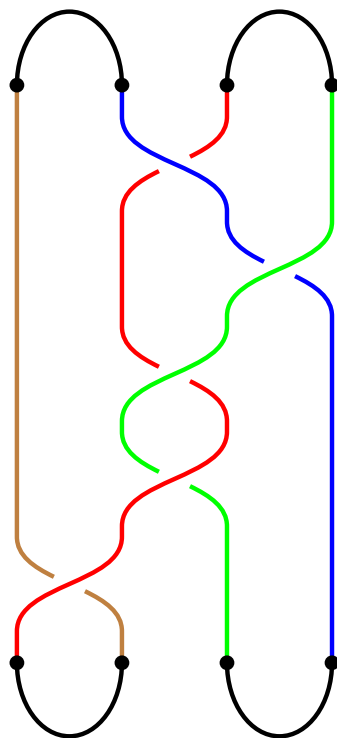
Topological quantum field theory (TQFT)



Put the system on a general surface

$$\Sigma \longmapsto V(\Sigma)$$

vector space of states



$$(M; K) \longmapsto Z(M; K)$$

Knot in a 3-manifold Knot invariant

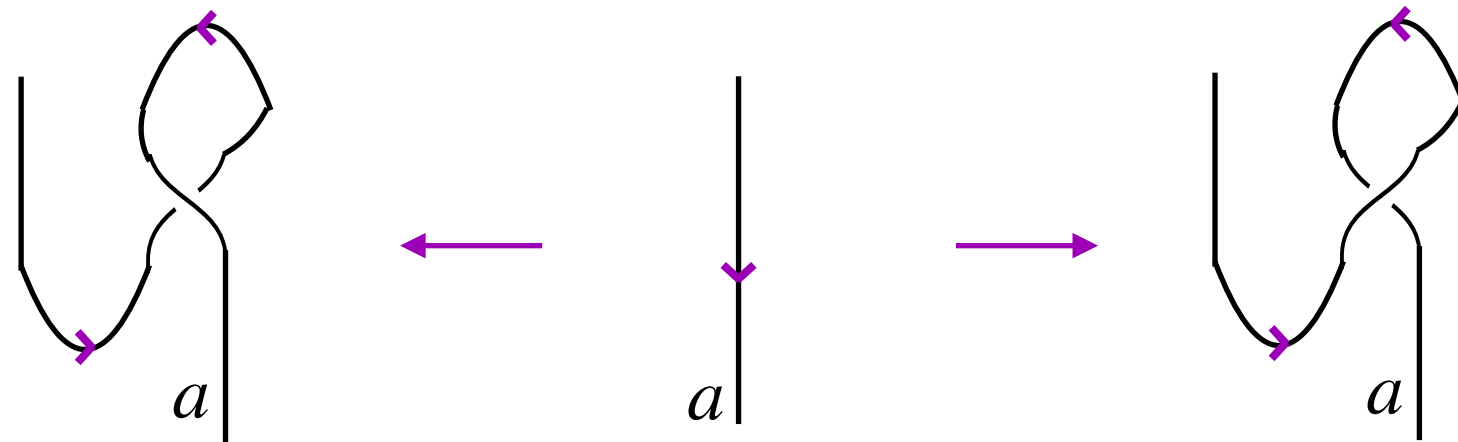
For an oriented knot/link K labelled by anyon a

$$\begin{array}{c} \text{Diagram 1: Knot } K \text{ with four segments labeled } a \\ \text{Diagram 2: Resolution of crossings with factor } (\dim_q a)^2 \\ \text{Diagram 3: Resulting diagram with labels } a, \bar{a} \text{ and a vertical line labeled } 1 \end{array}$$

- Can be computed using F - and R -symbols.
- Invariant under Reidemeister II and III moves.

$\langle K; a \rangle$ is an invariant of framed oriented links — Reshetikhin-Turaev invariant

- Reidemeister I move:



$\theta_a \langle K; a \rangle \quad \longleftarrow \quad \langle K; a \rangle \quad \longrightarrow \quad \overline{\theta}_a \langle K; a \rangle$

topological twist

Let $lk(K)$ be the linking number of K ,

$$J(K; a) := \theta_a^{-lk(K)} \langle K; a \rangle$$

is an invariant of un-framed oriented links

Fact: for the Fibonacci anyon, $J(K; \tau)$ is (a version of) the Jones polynomial of K evaluated at 5-th root of 1.

$\text{Vec}_{\mathbb{C}}$: category of vector spaces

objects: vector spaces, $V_1, V_2, \dots$

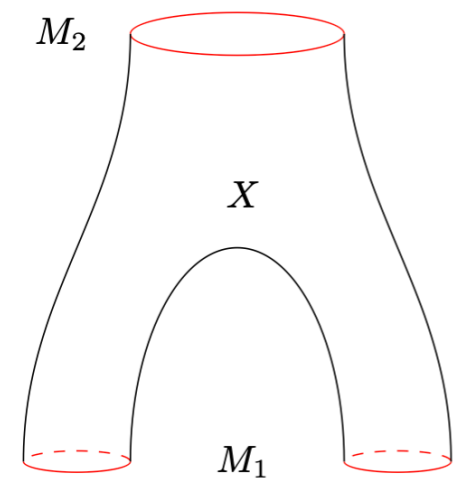
morphisms: linear maps $f : V_1 \rightarrow V_2$

$\text{Cob}(d + 1)$: category of cobordisms

objects: closed d -dim manifolds, $M_1, M_2, \dots$

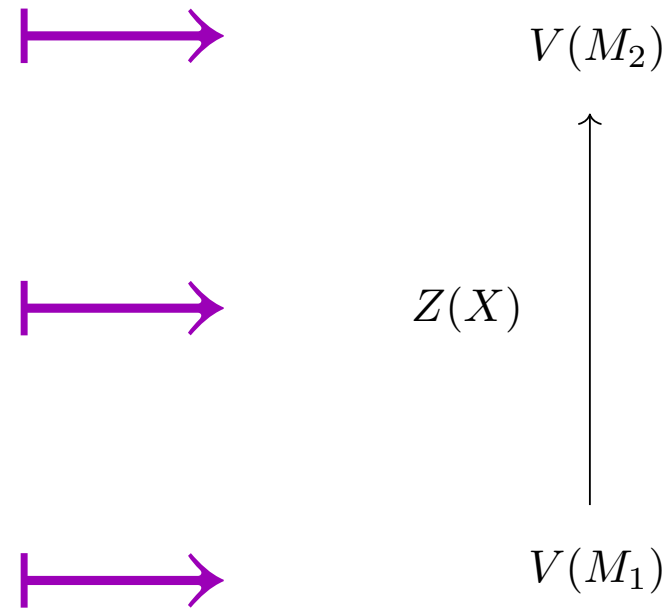
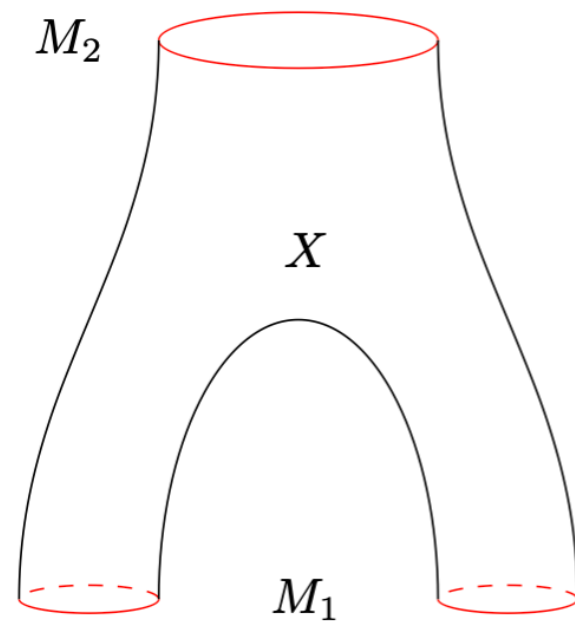
morphisms: $(d + 1)$ -dim cobordisms $X : M_1 \rightarrow M_2$

$$\text{i.e. } \partial X = M_1 \sqcup M_2$$

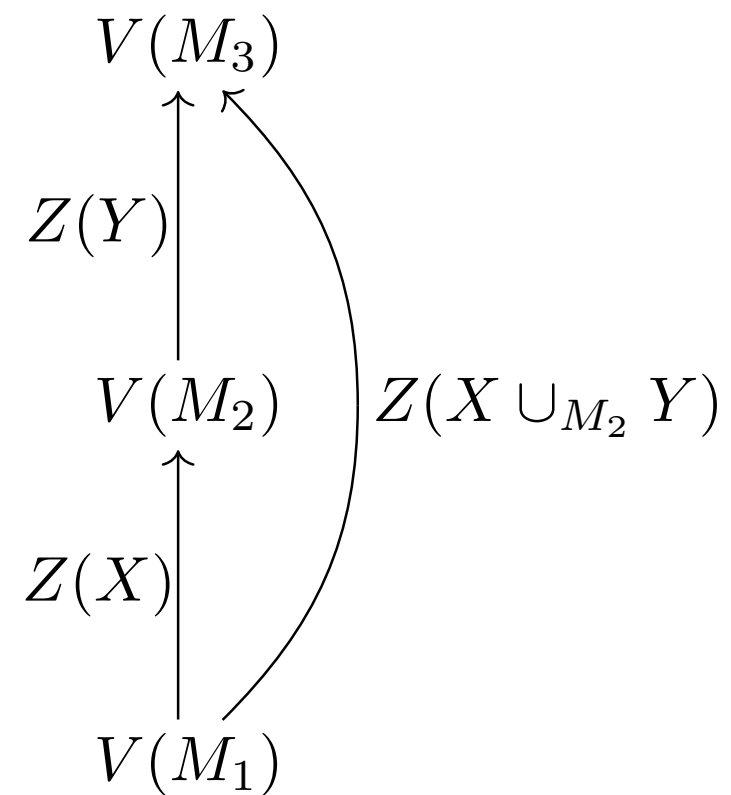
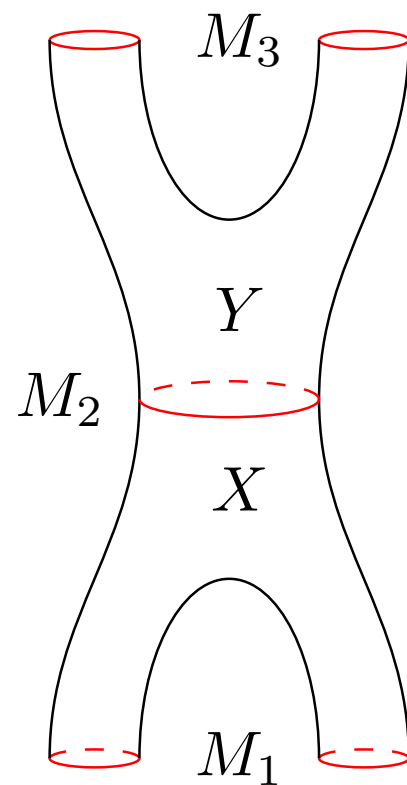


A $(d + 1)$ -dim topological quantum field theory (TQFT) is a symmetric functor

$$(V, Z) : \text{Cob}(d + 1) \rightarrow \text{Vec}_{\mathbb{C}}$$



Subject to (for instance)



- If X is closed, i.e., no boundary,

$$X : \emptyset \rightarrow \emptyset,$$

$$Z(X) : \mathbb{C} \rightarrow \mathbb{C}, \quad Z(X) \in \mathbb{C}$$

$Z(\cdot)$ is an invariant of closed $(d+1)$ -manifolds — quantum invariants

- Given $\phi : M \xrightarrow{\sim} M$, let M_ϕ be the mapping cylinder,

$$M_\phi : M \rightarrow M$$

$$Z(M_\phi) : V(M) \rightarrow V(M)$$

representation of the mapping class group: $\text{MCG}(M) \rightarrow \text{GL}(V(M))$

$$\phi \mapsto Z(M_\phi)$$

(2+1)-dim TQFTs

Modular tensor category $\mathcal{C}$ $\longrightarrow$ Witten-Reshetikhin-Turaev $Z_{\mathcal{C}}(\cdot)$

Let M be a 3-manifold and L be a surgery link for M

$$Z_{\mathcal{C}}(M) = \sum_{\text{anyon types } x \in \mathcal{C}} \dim_q(x) \left\langle \text{Reshetikhin-Turaev invariant} \right\rangle_x$$

weight

Volume Conjectures (Chen-Yang, '15)

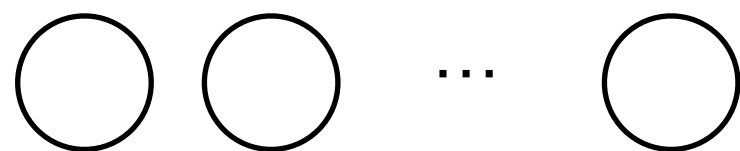
Let $\text{TLJ}(r)$ be the Temperley-Lieb-Jones category with the Kauffman variable $A = e^{\frac{\pi i}{r}}$, M be a closed oriented hyperbolic 3-manifold. Then for r odd,

$$\lim_{r \rightarrow \infty} \frac{4\pi}{r} \log |Z_{\text{TLJ}(r)}(M)| = \text{Vol}(M).$$

(3+1)-dimensions

Topological phases of matter:

- Anyons (point-like excitations) can only be fermions/bosons
- More interesting excitations are string-like. Motions of strings lead to representations of the motion groups of links.



N-component unlink



Loop braid group LB_n

$$\langle \sigma_1, \dots, \sigma_{n-1}, s_1, \dots, s_{n-1} \mid \dots \rangle$$

Algebraic theories: monoidal/braided 2-categories

Special cases: ribbon fusion categories, group-crossed braided fusion, spherical fusion 2-categories, etc.

(3+1)-TQFTs:

Construct TQFTs to distinguish smooth structures of 4-manifolds.

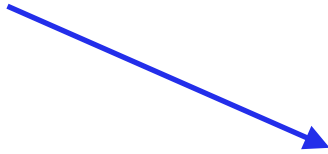
- Turaev-Viro type based on fusion 2-categories [Douglass-Reutter, '18]

No-go theorem:

Semisimple TQFTs cannot distinguish smooth structures [Reutter, '20]

Non-semisimple efforts:

- Trisection invariant based on Hopf algebras [Chaidez-Cotler-Cui, '19, '23]
- Skein TQFTs from ribbon categories [Costantino-Geer-Haïoun-Patureau-Mirand, '23]
- Invariants of 4D 2-handlebodies under 2-equivalence from ribbon cats [Beliakova-De Renzi-Faes, '25; Beliakova-De Renzi, '22]
- Khovanov homology and skein lasagna modules [Morrison-Walker-Wedrich, '22; Ren-Willis, '24]



can distinguish smooth structures

